Advanced Statistical Physics - Problem Set 2

Summer Term 2019

Due Date: Wednesday, April 17, 12:00 p.m., Hand in tasks marked with * to mailbox inside

ITP room 105b

Internet: Advanced Statistical Physics exercises

2. Field operators *

3 Points

The operators a_k^{\dagger} and a_k create or annihilate single particle states with momentum k, respectively. They obey the commutation relations $[a_k, a_{k'}]_{\zeta} = 0$, and $[a_k, a_{k'}^{\dagger}]_{\zeta} = \delta_{k,k'}$ with $\zeta = 1$ for bosons and $\zeta = -1$ for fermions. The field operator $\Psi(x)$ is defined as the Fourier transform

$$\Psi(x) = \frac{1}{\sqrt{L}} \sum_{k} a_k e^{ikx} .$$

Show that $\Psi(x)$ and its Hermitian conjugate $\Psi^{\dagger}(x)$ obey the commutation relations

$$[\Psi(x), \Psi^{\dagger}(y)]_{\zeta} = \delta(x-y)$$
.

3. Second quantization with field operators

4+5+1 Points

A many-particle state is described by the Hamiltonian

$$H = \sum_{j=1}^{N} \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_j^2} + U(x_j) \right) + \frac{1}{2} \sum_{i \neq j} V(x_i - x_j) ,$$

with U(x) being a scalar potential, V(x-x') being the two-particle interaction potential, and the wave function $\varphi_{\alpha}(x_1, x_2, \ldots, x_N)$ being an eigenstate of H with eigenvalue E_{α} . In the following, we denote the non-interacting part of the Hamiltonian by H_0 and the interaction part by H_{int} . In second quantized notation, the wave function φ_{α} is described by a state vector

$$|\varphi_{\alpha}\rangle = \int dx_1 dx_2 \dots dx_N \, \varphi_{\alpha}(x_1, x_2, \dots, x_N) \Psi^{\dagger}(x_1) \Psi^{\dagger}(x_2) \dots \Psi^{\dagger}(x_N) |0\rangle ,$$

with commutation relations of Ψ and Ψ^{\dagger} as in exercise 2, supplemented by the condition that the operator $\Psi(x)$ annihilates the vacuum state, $\Psi(x)|0\rangle = 0$. In second quantization, the Hamiltonian reads

$$H_s = \int dx \, \Psi^{\dagger}(x) \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + U(x) \right) \Psi(x) + \frac{1}{2} \int dx \, dy \, \Psi^{\dagger}(x) \Psi^{\dagger}(y) V(x-y) \Psi(y) \Psi(x) .$$

As before, we denote the non-interacting part by $H_{0,s}$ and the interaction part by $H_{\text{int},s}$. The aim of the task is to show that this is a consistent notation by proving

$$H_s|\varphi_{\alpha}\rangle = E_{\alpha}|\varphi_{\alpha}\rangle$$
.

(a)* Start with the non-interacting part $H_{0,s}$ and show that

$$H_{0,s}|\varphi_{\alpha}\rangle = \int dx_1 dx_2 \dots dx_N \left[H_0 \varphi_{\alpha}(x_1, x_2, \dots, x_N)\right] \Psi^{\dagger}(x_1) \Psi^{\dagger}(x_2) \dots \Psi^{\dagger}(x_N)|0\rangle$$
.

Hint: Use the result from exercise 2, i.e.

$$[\Psi(x), \Psi^{\dagger}(y)]_{\zeta} = \delta(x-y) ,$$

with $\zeta=\pm 1$ to cover bosons and fermions simultaneously.

(b) For the interaction part $H_{\text{int},s}$ show that

$$H_{\text{int},s}|\varphi_{\alpha}\rangle = \int dx_1 dx_2 \dots dx_N \left[H_{\text{int}}\varphi_{\alpha}(x_1, x_2, \dots, x_N)\right] \Psi^{\dagger}(x_1) \Psi^{\dagger}(x_2) \dots \Psi^{\dagger}(x_N)|0\rangle$$
.

(c) From your results of task (a) and (b) conclude that

$$H_s|\varphi_{\alpha}\rangle = E_{\alpha}|\varphi_{\alpha}\rangle$$
.