
Advanced Quantum Mechanics - Problem Set 4

Winter Term 2023/24

Due Date: Hand in solutions to problems marked with * as a single pdf file using Moodle before the lecture on **Thursday, 09.11.2023, 15:15**. The problem set will be discussed in the tutorials on Monday 13.11.2023 and Wednesday 15.11.2023.

Website: https://home.uni-leipzig.de/stp/Quantum_Mechanics_2_WS2324.html

Moodle: <https://moodle2.uni-leipzig.de/course/view.php?id=45746>

*1. SSH model

4+2+3 Points

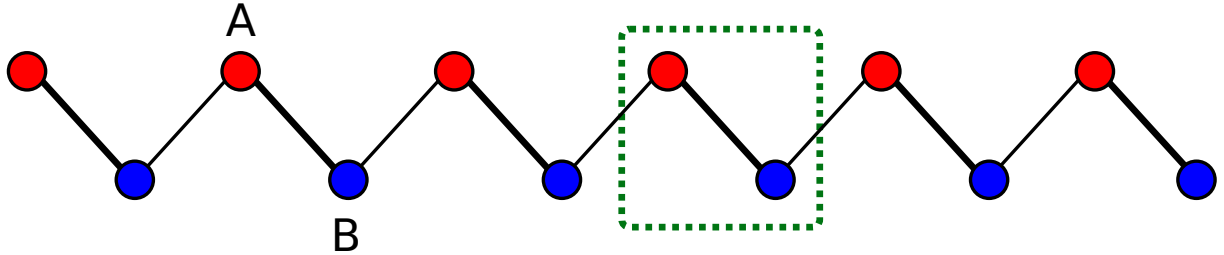


Figure 1: The SSH model. The red and blue circles symbolise different types of sites. The thin lines denote couplings with strength $t(1 - \delta)$ whilst the thick lines are couplings with strength $t(1 + \delta)$. The dashed square denotes a unit cell.

In this problem we consider the Su-Schrieffer-Heeger (SSH) model which describes spinless fermions hopping on a one-dimensional lattice with staggered hopping amplitudes (see the figure). The model contains two sub-lattices, A and B and has the following Hamiltonian

$$H = \sum_n t(1 + \delta)|n, A\rangle\langle n, B| + t(1 - \delta)|n + 1, A\rangle\langle n, B| + \text{h.c.}$$

Here h.c. stands for hermitian conjugate and $|n, A\rangle$ describes a state of site n , in sublattice A . t and δ are taken to be real parameters.

- (a) By Fourier transforming, $|n\rangle = \frac{1}{\sqrt{N}} \sum_k e^{-ink} |k\rangle$, show that the Hamiltonian can be written as $H(k) = \mathbf{d}(k) \cdot \boldsymbol{\sigma}$, where $\boldsymbol{\sigma}$ is the vector of Pauli matrices, and $d_x(k) = t(1 + \delta) + t(1 - \delta) \cos(k)$, $d_y(k) = t(1 - \delta) \sin(k)$, and $d_z(k) = 0$.

Hint: Write the wave function as a vector with two components describing the amplitudes on the A and B sublattices, respectively.

- (b) Calculate the energy eigenvalues of the system.
- (c) Plot your result from (b) for $\delta > 0$ and $\delta < 0$. What happens when $\delta = 0$?

2. Nearly free electron model

3+2+2+3 Points

Often it is sufficient to treat the periodic potential on a lattice as a small perturbation. For such problems it is useful to expand the periodic potential in a plane wave expansion which only contains waves with the periodicity of the reciprocal lattice, such that

$$U(\mathbf{x}) = \sum_{\mathbf{G}} U_{\mathbf{G}} e^{i\mathbf{G}\cdot\mathbf{x}},$$

where \mathbf{G} is a reciprocal lattice vector which satisfies $e^{i\mathbf{G}\cdot\mathbf{R}} = 1$, with \mathbf{R} denoting a point on the lattice. We moreover expand the wave functions in terms of a set of plane waves which satisfy the periodic boundary conditions of the problem

$$\psi(\mathbf{x}) = \sum_{\mathbf{k}} c_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{x}}.$$

- (a) Using the expansions above, show that the Schrödinger equation

$$\left[\frac{-\hbar^2 \nabla^2}{2m} + U(\mathbf{x}) \right] \psi(\mathbf{x}) = E\psi(\mathbf{x}),$$

can be written as

$$\left(\frac{\hbar^2 k^2}{2m} - E \right) c_{\mathbf{k}} + \sum_{\mathbf{G}} U_{\mathbf{G}} c_{\mathbf{k}-\mathbf{G}} = 0.$$

- (b) Perform the shift $\mathbf{q} = \mathbf{k} + \mathbf{K}$, where \mathbf{K} is a reciprocal lattice vector which ensures that we can always find a \mathbf{q} which lies in the first Brillouin zone¹, and show that the Schrödinger equation now gives

$$\left(\frac{\hbar^2}{2m} (\mathbf{q} - \mathbf{K})^2 - E \right) c_{\mathbf{q}-\mathbf{K}} + \sum_{\mathbf{G}} U_{\mathbf{G}-\mathbf{K}} c_{\mathbf{q}-\mathbf{G}} = 0.$$

- (c) Consider for concreteness a one-dimensional chain, but in the simple case where only the leading Fourier component contributes to the potential

$$U(x) = 2U_0 \cos \frac{2\pi x}{a}.$$

Explain how your result in (b) can be used to calculate the energy of the system.

- (d) Suppose now that U_0 is very small. Near $q = \pi/a$ the Schrödinger equation reduces to

$$\begin{pmatrix} \frac{\hbar^2}{2m} \left(q - \frac{2\pi}{a} \right)^2 - E & U_0 \\ U_0 & \frac{\hbar^2 q^2}{2m} - E \end{pmatrix} \begin{pmatrix} c_1 \\ c_0 \end{pmatrix} = 0.$$

Calculate and plot the energy eigenvalues. What happens at $q = \pi/a$?

¹As an example of a Brillouin zone consider the simple cubic lattice with sides of length a . The lattice vectors can be written as $\mathbf{R}_1 = a\hat{x}$, $\mathbf{R}_2 = a\hat{y}$, and $\mathbf{R}_3 = a\hat{z}$. In reciprocal space the basis vectors become $\mathbf{b}_1 = \frac{2\pi}{a}\hat{x}$, $\mathbf{b}_2 = \frac{2\pi}{a}\hat{y}$, and $\mathbf{b}_3 = \frac{2\pi}{a}\hat{z}$. In this case the first Brillouin zone is the region $-\pi/a \leq k_i < \pi/a$ (where $i = x, y, z$). The reciprocal lattice vectors can be written as $\mathbf{K} = \sum_i n_i \mathbf{b}_i$ (where $n_i \in \mathbb{Z}$). Therefore, for arbitrary \mathbf{k} it is possible to find $\mathbf{q} = \mathbf{k} + \mathbf{K}$ so that \mathbf{q} lies in the first Brillouin zone.